Theory of Equivalent Networks and Some of its Applications

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#### **1** Introduction

We consider the following network  $\Omega$ . The nodes of  $\Omega$  are the points in 3D space. The set of nodes is denoted by  $\Omega_n$ . We define some subset  $\Omega_b$  of  $\Omega_n$  as *boundary* nodes, so that cardinality of  $\Omega_b$  N<sub>b</sub> =  $card\Omega_b \geq 2$ . All other nodes of  $\Omega_n$  are called *interior*. The set of interior nodes is denoted by  $\Omega_i$ . N<sub>i</sub> =  $card\Omega_i \geq 0$ .

 $\begin{aligned} \Omega_n &= \Omega_b \cup \Omega_i. \\ \Omega_b \cap \Omega_i &= \emptyset. \\ \mathbf{N} &= card\Omega_n = \mathbf{N}_b + \mathbf{N}_i. \end{aligned}$ 

For two different nodes p and q belonging to  $\Omega_n$  the number  $\gamma(pq)$  is called *conductance* of pq. The function  $\gamma$ , called *conductivity*, is a non-negative real function. Two nodes p and q are called *neighbors* if  $\gamma(pq) > 0$ . The set of all neighbors of node p is called N(p). If p and q are neighbors pq is called an *edge* or a *conductor*.

A *network of conductors* is a network with defined set of boundary nodes and conductivity function. Below we will call a network of conductors just a network.

For a node p belonging to  $\Omega_n u(p)$  denotes the potential at p. A function  $u: \Omega_n \to \mathbb{R}$  gives the *current* from node p to  $q \in N(p)$  by

Ohm's Law

$$I(pq) = \gamma(pq)(u(p) - u(q)).$$

For  $q \notin N(p)$ 

$$I(pq) = I(qp) = 0.$$

For a boundary node q

$$I(q) = \sum_{p \in N(q)} I(qp)$$

is called a *boundary current*.

A function u is called  $\gamma$ -harmonic function, if for each interior node p

$$\sum_{q \in N(p)} I(pq) = 0.$$

According to *Kirchoff's Law*, if a function  $\phi$  is defined at the boundary nodes, the network  $\Omega$  will acquire a unique  $\gamma$ -harmonic function u with  $u(p) = \phi(p)$  for all boundary nodes p.

For a network  $\Omega$  with a conductivity function on it the linear map  $\Lambda: \phi \to I_{\phi}$  is defined by  $\Lambda \phi = I_{\phi}$ , where  $\phi$  is a vector of boundary potentials  $I_{\phi}$  is the corresponding vector of boundary currents. This map is called Dirichlet-to-Neumann map. ¿From the maximum principle and Kirchoff's and Ohm's Laws it can be proved that any  $\Lambda$  has following properties (See [1])

- it is symmetric
- all non-diagonal entries are non-positive
- total of all entries in each row equals 0.

## 2 Plane equivalent of a network

**Definition 2.1** The *Kirchoff matrix* of the network  $\Omega$  is a square NxN matrix  $K(\Omega) = \{k_{ij}\}$  where

$$k_{ij} = -\gamma(p_i p_j)$$
 if  $i \neq j$ 

and

$$k_{ii} = \sum_{q \in N(p_i)} \gamma(p_i q).$$

The matrix  $K(\Omega)$  contains all information about conductivity function of network  $\Omega$ . We denote the set of Kirchoff matrices by  $\mathcal{K}$ . A matrix  $M = \{m_{ij}\}$  belongs to  $\mathcal{K}$  if

$$m_{ij} = m_{ji} \le 0 \qquad \text{if} \qquad i \ne j,$$
  

$$m_{ii} > 0,$$
  

$$\sum_{j=1}^{N} m_{ij} = 0.$$

By this definition the set of Dirichlet-to-Neumann maps belongs to  $\mathcal{K}$ .

If  $\Omega$  is a network with no interior nodes, by Kirchoff's and Ohm's Laws

$$K(\bar{\Omega})\phi = I_{\phi},$$

therefore,

$$K(\bar{\Omega}) = \Lambda(\bar{\Omega}).$$

This proves the next statements.

**Statement 2.2** The set of Dirichlet-to-Neumann maps *equals*  $\mathcal{K}$ .

Statement 2.3 For any  $\Lambda \in \mathcal{K}$  there exists unique network  $\Omega$  with no interior nodes such that  $\Lambda$  represents the Dirichlet-to-Neumann map for  $\Omega$ .

**Statement 2.4** The map from networks with no interior nodes to Dirichlet-to-Neumann maps is  $1 \leftrightarrow 1$ .

**Definition 2.5** A plane equivalent of network  $\Omega$  is a network  $\overline{\Omega}$  with no interior nodes which has the same Dirichlet-to-Neumann map as  $\Omega$ .

**Theorem 2.6** Any resistor network has unique plane equivalent.

**Definition 2.7** Two networks are *equivalent* if they have the same plane equivalent.

For a network  $\Omega$  with a conductivity function on it the linear map  $\Phi: \phi \to u_{\phi}$  is defined by  $\Phi \phi = u_{\phi}$ , where  $\phi$  is a vector of boundary potentials and  $u_{\phi}$  is the corresponding vector of potentials on nodes of  $\Omega$ .  $\Phi$  is represented by an NxN<sub>b</sub> matrix.

We number nodes of a network  $\Omega$  so that boundary nodes go first. That is,

$$p_i \in \Omega_b \Leftrightarrow i \leq \mathcal{N}_b$$

and

$$p_i \in \Omega_i \Leftrightarrow i > \mathcal{N}_b$$

We write  $K(\Omega)$  and  $\Phi(\Omega)$  in block form,

$$\Phi(\Omega) = \begin{bmatrix} I \\ \Phi_{\Omega} \end{bmatrix} \text{ and } K(\Omega) = \begin{bmatrix} K' & M^T \\ M & \Xi \end{bmatrix}$$

where I is the identity matrix. Now using Kirchoff's and Ohm's Laws we obtain that

$$\begin{array}{c|cccc}
K' & M^T \\
\hline
M & \Xi \\
\end{array} & \hline
\Phi_{\Omega} \\
\end{array} = \boxed{\Lambda(\Omega)} \\
0
\end{array}$$
(1)

where  $\theta$  is a zero matrix. Since  $\Xi$  is non-singular, we can obtain that

$$\Lambda(\Omega) = K(\bar{\Omega}) = K' - M^T \Xi^{-1} M.$$
(2)

This equation expresses the Dirichlet-to-Neumann map of a network in terms of blocks of this network's Kirchoff matrix.

### **3** Blocks in Dirichlet-to-Neumann maps

We consider a network  $\Omega$  with Dirichlet-to-Neumann map  $\Lambda$ . We write the Kirchoff matrix of  $\Omega$  in block form,

$$K(\Omega) = \begin{array}{c|c} K' & M^T \\ \hline M & \Xi \end{array}$$

Let A be a n by n square submatrix of K' corresponding to rows  $i_1, ..., i_n$  and columns  $j_1, ..., j_n$ . Let C be the  $N_i \ge n$  matrix formed by choosing columns  $j_1, ..., j_n$  of block M. Let B be the  $n \ge N_i$  matrix

formed by choosing rows  $i_1, ..., i_n$  of block  $M^T$ .  $\Lambda_A$  is the submatrix of  $\Lambda$  corresponding to rows  $i_1, ..., i_n$  and columns  $j_1, ..., j_n$ .

By formula (2)

$$\Lambda_A = A - B \Xi^{-1} C.$$

Now we have that

A	В	Ι	0	=	$\Lambda_A$	В
С	[I]	$-\Xi^{-1}C$	Ι		0	[1]

If

$$W = \begin{bmatrix} A & B \\ C & \Xi \end{bmatrix}$$

we obtain that

$$|W| = |\Lambda_A||\Xi|. \tag{3}$$

Taking A to be an entry of K' we can find any entry of  $\Lambda$  in terms of two determinants of blocks of  $K(\Omega)$ .

$$\lambda_{ij} = \begin{vmatrix} k_{ij}' & m_i^T \\ m_j & \Xi \end{vmatrix} / |\Xi|.$$

Since  $\Xi$  is positive definite we have that

$$sign|\Lambda_A| = sign|W|.$$

# 4 Theorem about elimination of a subnetwork

**Definition 4.1** A network  $\hat{\Omega} = {\hat{\Omega}_n, \hat{\Omega}_b, \hat{\gamma}}$  is a *subnetwork* of  $\Omega = {\Omega_n, \Omega_b, \gamma}$  if

 $\hat{\Omega}_i \subset \Omega_i.$ 

$$\hat{\Omega}_b = (\bigcup_{p \in \hat{\Omega}_i} N(p)) \setminus \hat{\Omega}_i.$$
  
$$p, q \in \hat{\Omega}_n \Rightarrow \hat{\gamma}(pq) = \gamma(pq).$$

Any subnetwork is defined by its set of interior nodes. The Kirchoff matrix  $\hat{K}$  of  $\hat{\Omega}$  is a principal submatrix of K.

**Definition 4.2** *Elimination* of a subnetwork  $\hat{\Omega}$  from a network  $\Omega$  is a transformation of  $\Omega$  to  $\Psi$  denoted by  $\Omega \ominus \hat{\Omega}$  if

$\Psi_b = \Omega_b$		
$\Psi_i = \Omega_i \backslash \hat{\Omega}_i.$		
$\gamma_{\Psi}(pq) = \gamma_{\Omega}(pq)$	if	$p \notin \hat{\Omega}_n$ .
$\gamma_{\Psi}(pq) = \gamma_{\hat{\bar{\Omega}}}(pq)$	if	$p,q\in\hat{\Omega}_b.$

where  $\gamma_{\hat{\Omega}}$  is the conductivity function on the plane equivalent of  $\hat{\Omega}$ . In other words to eliminate some subnetwork from a network is to replace this subnetwork with its plane equivalent.

We number the nodes so that the boundary nodes of  $\Omega$  go first and interior nodes of  $\hat{\Omega}$  go last. That is,

 $p_i \in \Omega_b \Leftrightarrow i \leq card\Omega_b$ 

and

$$p_i \in \hat{\Omega}_i \Leftrightarrow i > card\Omega - card\hat{\Omega}_i$$

Now we have that

	K'	$P^T$	$U^T$
$K(\Omega) =$	Р	S	$R^T$
	U	R	C

Taking into account that C is non-singular and using Definition 4.2 and formula (2) we can obtain that if  $\Psi = \Omega \ominus \hat{\Omega}$  then

$$K(\Psi) = \begin{bmatrix} K' - U^T C^{-1} U & P^T - U^T C^{-1} R \\ P - R^T C^{-1} U & S - R^T C^{-1} R \end{bmatrix} = \Lambda(\{\Psi_n, \Psi_n, \gamma_{\Psi}\})$$

We define linear maps  $\Phi(\Omega)$  and  $\Phi(\Psi)$  as it was done in Section 2. Solving equation (1) we obtain that

$$\Phi_{\Psi} = -(S - R^T C^{-1} R)^{-1} (P - R^T C^{-1} U)$$

and

$$\Phi(\Omega) = \begin{bmatrix} I \\ \Phi_{\Psi} \\ -C^{-1}(T + R\Phi_{\Psi}) \end{bmatrix} = \begin{bmatrix} \Phi(\Psi) \\ -C^{-1}(T + R\Phi_{\Psi}) \end{bmatrix}$$

The non-singularity of  $S - R^T C^{-1} R$  follows from (3). Now it is easy to show that

$$\Lambda(\Psi) = \Lambda(\Omega).$$

We proved the following theorem.

**Theorem 4.4** Suppose  $\Psi = \Omega \ominus \hat{\Omega}$  where  $\hat{\Omega}$  is a subnetwork of  $\Omega$ . Then  $\Psi$  and  $\Omega$  are equivalent. That is  $\Lambda(\Psi) = \Lambda(\Omega)$ . The matrix  $\Phi(\Psi)$  is an upper part of  $\Phi(\Omega)$ . Therefore,

$$I_{\Psi}(pq) = I_{\Omega}(pq)$$
 if  $p, q \notin \Psi_N$ .

In other words, elimination of a subnetwork does not affect potentials and currents outside of it.

## 5 Recoverable networks

We will say that we know the *shape* of a network if we know all neighbors for each node of this network.

**Definition 5.1** A network is *recoverable* if there does not exist a network with the same shape, same plane equivalent (or Dirichlet-to-Neumann map), and different conductivity function.

A network is *non-recoverable* if it is not recoverable.

By Definition 5.1 two recoverable networks with the same shape have different plane equivalents.

**Statement 5.2** All subnetworks of a recoverable network are recoverable.

Elimination of node p is an elimination of a subnetwork  $\hat{\Omega}$  such that  $\hat{\Omega}_i = p$ .

Elimination of a subnetwork  $\hat{\Omega}$  from  $\Omega$  results in the same network as elimination of all interior nodes of  $\hat{\Omega}$  in any order. Elimination of some set S of interior nodes of  $\Omega$  results in the same network as elimination of the subnetwork  $\hat{\Omega}$  if  $\hat{\Omega}_i = S$ .

By Theorem 4.4 elimination of an interior node p of a network  $\Omega$ does not affect plane equivalent of any subnetwork  $\hat{\Omega}$  of  $\Omega$  if  $p \in \hat{\Omega}_i$ .

## 6 Chain elimination

We consider a network  $\Omega = \{\Omega_n, \Omega_b, \gamma\}$  with  $\Omega_b \neq \Omega_n$ . We will eliminate interior nodes of  $\Omega$  one by one. The network obtained after  $j^{th}$  elimination is called  $\Omega^j$ . Let

$$E^j = \Omega_n \backslash \Omega_n^j.$$

Let  $\hat{\Omega}^j$  be the subnetwork of  $\Omega$  with  $\hat{\Omega}^j_i = E^j$ . Then

$$\Omega^j = \Omega \ominus \hat{\Omega}^j.$$

By Theorem 4.4

 $\Lambda(\Omega) = \Lambda(\Omega^j).$ 

Let  $\Psi^j = \{\Omega_n, \Psi^j_b, \gamma\}$  where

$$\Psi_h^j = \Omega_n \backslash E^j.$$

 $\Psi^{N_i} = \Omega$ . By formula (2) we obtain that

$$K(\Omega \ominus \hat{\Omega}^{j}) = K(\Omega^{j}) = \Lambda(\Psi^{j}).$$
$$K(\Omega) = \Lambda(\Psi^{0}).$$

# 7 Specific comments

Formula (2) gives a direct algorithm for finding the Dirichlet-to-Neumann map for any network with known conductivity function. For networks with two boundary nodes this formula represents the generalized parallel and series laws of electrical circuit theory.