

# BASIC VON NEUMANN ALGEBRA THEORY

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## 1. INTRODUCTION

The theory of von Neumann algebras and von Neumann dimensions allows one to measure some infinite-dimensional subspaces in a Hilbert space, by assigning to them a notion of “dimension” (not necessarily an integer). We quickly review some very basic notions in the theory. We follow the presentation in [Shu93]. See also [Lüc02] and [Pan96] for proofs and a more thorough treatment.

## 2. VON NEUMANN ALGEBRAS AND FACTORS

Let  $H$  be a Hilbert space over  $\mathbb{C}$  with the (Hermitian) inner product  $\langle \cdot, \cdot \rangle$ .

Let  $\mathcal{B}(H)$  be the algebra of all bounded linear operators on  $H$ . This is Banach space with respect to the operator norm. Moreover, this is a  $C^*$ -algebra:

- (i)  $\mathcal{B}(H)$  is a Banach algebra, meaning  $\|AB\| \leq \|A\|\|B\|$  for all  $A, B \in \mathcal{B}(H)$ .
- (ii)  $\mathcal{B}(H)$  is a  $*$ -algebra, i.e. it is closed under the operation of taking adjoints of operators, where  $A \mapsto A^*$  is defined by the usual identity  $\langle Ax, y \rangle = \langle x, A^*y \rangle$  for all  $x, y \in H$ .
- (iii)  $(AB)^* = B^*A^*$  for all  $A, B \in \mathcal{B}(H)$ .
- (iv)  $\|A^*A\| = \|A\|^2$  for all  $A \in \mathcal{B}(H)$ .

For any subset  $\mathcal{M} \subseteq \mathcal{B}(H)$ , its *commutant* is the subalgebra

$$\mathcal{M}' = \{A \in \mathcal{B}(H) : AB = BA, \forall B \in \mathcal{M}\}$$

in  $\mathcal{B}(H)$ . Clearly, the identity operator  $I$  is in  $\mathcal{M}'$ .

**Definition 2.1.** A *von Neumann algebra* on  $H$  is a subalgebra  $\mathcal{M} \subseteq \mathcal{B}(H)$  such that

- (i)  $\mathcal{M}$  is a  $*$ -subalgebra.
- (ii)  $\mathcal{M}'' = \mathcal{M}$ .

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It follows from  $\mathcal{M}'' = \mathcal{M}$  that  $I \in \mathcal{M}$ .

*Remark 2.2.* In Definition 2.1, the slightly unmotivated condition  $\mathcal{M}'' = \mathcal{M}$  can be replaced by either of the following two more natural conditions:

- (a)  $I \in \mathcal{M}$ , and  $\mathcal{M}$  is closed in the weak operator topology.
- (b)  $I \in \mathcal{M}$ , and  $\mathcal{M}$  is closed in the strong operator topology.

This is usually called the *von Neumann Double Commutant Theorem*.

**Definition 2.3.** Let  $\mathcal{M}$  be a von Neumann algebra on  $H$ .

- (i)  $\mathcal{M}$  is called a *factor* if its center is trivial, i.e.  $\mathcal{M} \cap \mathcal{M}' = \{\lambda I : \lambda \in \mathbb{C}\}$ .
- (ii) The set of all *orthogonal projections* in  $\mathcal{M}$  is denoted by  $\text{Pr}(\mathcal{M})$ . In other words

$$\text{Pr}(\mathcal{M}) = \{A \in \mathcal{M} : A^2 = A = A^*\}.$$

Factors can be thought of as “indecomposable” von Neumann algebras: every von Neumann algebra on a (separable) Hilbert space is isomorphic to a direct integral of factors. For our purpose, their significance lies in the fact that if  $\mathcal{M}$  is a factor, then there is *at most one trace function* (up to a constant factor) on  $\mathcal{M}$ .

### 3. VON NEUMANN TRACE

**Definition 3.1.** Let  $\mathcal{M}$  be a von Neumann algebra on  $H$ . Let

$$\mathcal{M}^+ = \{A \in \mathcal{M} : A \geq 0\}.$$

(Here  $A \geq 0$  means  $\langle Ax, x \rangle \in \mathbb{R}^{\geq 0}$ ,  $\forall x \in H$ ).

A *trace* on  $\mathcal{M}$  is a function  $\tau : \mathcal{M}^+ \rightarrow [0, +\infty]$  satisfying:

- (Linear)  $\tau(aA + bB) = a\tau(A) + b\tau(B)$ ,  $\forall a, b \in \mathbb{R}^{\geq 0}$  and  $\forall A, B \in \mathcal{M}^+$ .
- (Tracial)  $\tau(A^*A) = \tau(AA^*)$ ,  $\forall A \in \mathcal{M}$ .
- (Faithful)  $\tau(A) = 0 \Rightarrow A = 0$ .
- (Normal) If  $\{A_\alpha\}$  is an increasing net on  $\mathcal{M}^+$  converging (in operator norm topology) to  $A$ , then  $\{\tau(A_\alpha)\}$  is an increasing net on  $\mathbb{R}^{\geq 0}$  converging to  $\tau(A)$ .
- (Semifinite)  $\tau(A) = \sup\{\tau(B) : B \in \mathcal{M}^+, \tau(B) < +\infty, A - B \geq 0\}$ ,  $\forall A \in \mathcal{M}^+$ .

The most basic example is the following:

**Example 3.2.** Clearly,  $\mathcal{M} = \mathcal{B}(H)$  is a von Neumann algebra. Moreover, it is a factor, hence there is at most one trace function  $\text{Tr}$  (up to scaling) associated to it. Let  $\{u_\alpha\}_{\alpha \in J}$  be an orthonormal basis for  $H$ . For  $A \in \mathcal{B}(H)^+$  define:

$$(1) \quad \text{Tr}(A) = \sum_{\alpha \in J} \langle Au_\alpha, u_\alpha \rangle,$$

where the sum is understood as the supremum of finite partial sums. Then  $\text{Tr}$  defines a trace function according to Definition 3.1, and it is independent of the choice of the orthonormal basis. Any other trace on  $\mathcal{B}(H)$  is of the form  $c \text{Tr}$  for some  $c > 0$ .

4. VON NEUMANN DIMENSION

Now we can describe the beautiful approach to assigning a notion of “dimension” to certain infinite dimensional subspaces of  $H$ .

**Definition 4.1.** Let  $\mathcal{M}$  be a von Neumann algebra on  $H$ .

- (i) A subspace  $L \subseteq H$  is *admissible* with respect to  $\mathcal{M}$  (or  $\mathcal{M}$ -admissible) if there is an orthogonal projection  $A \in \text{Pr}(\mathcal{M})$  such that  $L = \text{Image}(A)$ . The set of all  $\mathcal{M}$ -admissible subspaces of  $H$  is the  *$\mathcal{M}$ -admissible Grassmannian*:

$$\mathbf{Gr}(H, \mathcal{M}) = \{L \subseteq H : L = \text{Image}(A) \text{ for some } A \in \text{Pr}(\mathcal{M})\}.$$

- (ii) Given a trace  $\tau$  on  $\mathcal{M}$ , define the corresponding *dimension* function

$$\dim_\tau : \mathbf{Gr}(H, \mathcal{M}) \rightarrow [0, +\infty]$$

by

$$L = \text{Image}(A) \mapsto \tau(A).$$

It is elementary to check that this notion of dimension generalizes the dimension theory for finite-dimensional vector spaces.

The following (expected) properties of  $\dim_\tau$  are easy to verify:

- (i)  $\dim_\tau(L) = 0 \iff L = \{0\}$ .
- (ii)  $L_1 \subseteq L_2 \Rightarrow \dim_\tau(L_1) \leq \dim_\tau(L_2)$ .
- (iii)  $L_1 \perp L_2 \Rightarrow \dim_\tau(L_1 \oplus L_2) = \dim_\tau(L_1) + \dim_\tau(L_2)$ .
- (iv) If  $U \in \mathcal{M}$  is *unitary* (i.e.  $U^{-1} = U^*$ ), then  $\dim_\tau(U(L)) = \dim_\tau(L)$ .

5. TENSOR PRODUCTS

Let  $(H_1, \langle \cdot, \cdot \rangle_1)$  and  $(H_2, \langle \cdot, \cdot \rangle_2)$  be two Hilbert spaces. Let  $H_1 \otimes_{\text{alg}} H_2$  denote the algebraic (vector space) tensor product. The vector space  $H_1 \otimes_{\text{alg}} H_2$  is endowed with an inner product defined by  $\langle x_1 \otimes x_2, y_1 \otimes y_2 \rangle = \langle x_1, y_1 \rangle \langle x_2, y_2 \rangle$ . The Hilbert space tensor product  $H_1 \otimes H_2$  is, by definition, the (metric space) completion of  $H_1 \otimes_{\text{alg}} H_2$  with respect to the inner product.

If  $A_1 \in \mathcal{B}(H_1)$  and  $A_2 \in \mathcal{B}(H_2)$ , we can form the algebraic tensor product  $A_1 \otimes_{\text{alg}} A_2 : H_1 \otimes_{\text{alg}} H_2 \rightarrow H_1 \otimes_{\text{alg}} H_2$ . It can be extended (by continuity) to an operator  $A_1 \otimes A_2 \in \mathcal{B}(H_1 \otimes H_2)$ .

If  $\mathcal{M}_1 \subseteq \mathcal{B}(H_1)$  and  $\mathcal{M}_2 \subseteq \mathcal{B}(H_2)$  are von Neumann algebras we may form the von Neumann tensor product  $\mathcal{M}_1 \otimes \mathcal{M}_2$ . This is the strong (or weak) completion of the set of all finite linear combinations of tensor products  $A_1 \otimes A_2$  where  $A_1 \in \mathcal{M}_1$  and  $A_2 \in \mathcal{M}_2$ . Equivalently,  $\mathcal{M}_1 \otimes \mathcal{M}_2$  is the von Neumann algebra generated by the algebraic tensor product  $\mathcal{M}_1 \otimes_{\text{alg}} \mathcal{M}_2$  inside  $\mathcal{B}(H_1) \otimes \mathcal{B}(H_2)$ .

6. VON NEUMANN ALGEBRAS ASSOCIATED WITH A DISCRETE GROUP

We will now focus on the von Neumann algebras and dimensions that arise in the presence of group actions.

**6.1. Group von Neumann algebras.** Let  $\mathbf{G}$  be a discrete group (topological group equipped with the discrete topology). Consider the Hilbert space:

$$\ell^2(\mathbf{G}) = \{f \mid f: \mathbf{G} \rightarrow \mathbb{C}, \|f\|^2 = \sum_{g \in \mathbf{G}} |f(g)|^2 < +\infty\},$$

where, as usual, the sum is understood as the supremum of finite partial sums. An orthonormal basis is given by

$$\{\delta_g: g \in \mathbf{G}\},$$

where  $\delta_g(h) = 1$  if  $g = h$  and 0 otherwise.

Translations operators generate (two) natural von Neumann algebra on  $\ell^2(\mathbf{G})$ , which we now describe. For each  $g \in \mathbf{G}$ , we have the following left and right translation operators  $L_g, R_g \in \mathcal{B}(\ell^2(\mathbf{G}))$ :

$$L_g(f)(h) = f(g^{-1}h) \quad , \quad h \in \mathbf{G},$$

$$R_g(f)(h) = f(hg) \quad , \quad h \in \mathbf{G}.$$

In terms of the orthonormal basis  $\{\delta_g: g \in \mathbf{G}\}$  we have

$$L_g(\delta_h) = \delta_{gh} \quad , \quad R_g(\delta_h) = \delta_{hg^{-1}}.$$

Clearly, both  $L_g$  and  $R_g$  are unitary:

$$(L_g)^* = (L_g)^{-1} = L_{g^{-1}} \quad , \quad (R_g)^* = (R_g)^{-1} = R_{g^{-1}}.$$

It is easy to see

$$L_{g_1 g_2} = L_{g_1} L_{g_2} \quad , \quad R_{g_1 g_2} = R_{g_1} R_{g_2}.$$

In fact, the assignments  $g \mapsto L_g$  and  $g \mapsto R_g$  are the usual left and right regular representations of  $\mathbf{G}$ .

Now we are ready to define the relevant group von Neumann algebras:

**Definition 6.1.**

- (i)  $\mathcal{M}_l = \mathcal{M}_l(\mathbf{G})$  is the von Neumann algebra generated by  $\{L_g: g \in \mathbf{G}\} \subseteq \mathcal{B}(\ell^2(\mathbf{G}))$ .
- (ii)  $\mathcal{M}_r = \mathcal{M}_r(\mathbf{G})$  is the von Neumann algebra generated by  $\{R_g: g \in \mathbf{G}\} \subseteq \mathcal{B}(\ell^2(\mathbf{G}))$ .

It turns out

$$\mathcal{M}_l = (\mathcal{M}_r)' \quad , \quad \mathcal{M}_r = (\mathcal{M}_l)'$$

so one might alternatively define  $\mathcal{M}_l$  (respectively  $\mathcal{M}_r$ ) as the algebra of  $\mathbf{G}$ -equivariant bounded operators on  $\ell^2(\mathbf{G})$ .

Also,  $\mathcal{M}_l$  and  $\mathcal{M}_r$  are factors if and only if all non-trivial conjugacy classes in  $\mathbf{G}$  are infinite. For example, this is the case when  $G$  is a free group on at least two generators.

Finally, a trace on  $\mathcal{M}_r$  (and also on  $\mathcal{M}_l$ ) is defined by

$$(2) \quad \tau(A) = \langle A\delta_h, \delta_h \rangle,$$

for  $h \in \mathbf{G}$ . This is independent of the choice of  $h$ , so it is convenient to use the  $h = \text{id}$ , the group identity.

**6.2. Hilbert  $\mathbf{G}$ -modules, and  $\mathbf{G}$ -dimensions.** By a Hilbert  $\mathbf{G}$ -module we mean a Hilbert space together with a (left) unitary action of the discrete group  $\mathbf{G}$ .

**Definition 6.2.** A free Hilbert  $\mathbf{G}$ -module is a Hilbert  $\mathbf{G}$ -module which is unitarily isomorphic to

$$\ell^2(\mathbf{G}) \otimes H,$$

where  $H$  is a Hilbert space with the trivial  $\mathbf{G}$ -action and the action of  $\mathbf{G}$  on  $\ell^2(\mathbf{G})$  is by left translations. In other words, the representation of  $\mathbf{G}$  on  $\ell^2(\mathbf{G}) \otimes H$  is given by  $g \mapsto L_g \otimes I$ .

Let  $\{u_\alpha\}_{\alpha \in J}$  be an orthonormal basis for  $H$ . Then we have an orthogonal decomposition

$$\ell^2(\mathbf{G}) \otimes H = \bigoplus_{\alpha \in J} \ell^2(\mathbf{G})^{(\alpha)},$$

where  $\ell^2(\mathbf{G})^{(\alpha)} = \ell^2(\mathbf{G}) \otimes u_\alpha$  is just a copy of  $\ell^2(\mathbf{G})$ . This means

$$\ell^2(\mathbf{G}) \otimes H = \left\{ \sum_{\alpha \in J} f_\alpha \otimes u_\alpha : f_\alpha \in \ell^2(\mathbf{G}), \sum_{\alpha \in J} \|f_\alpha\|^2 < +\infty \right\}.$$

Moreover, an orthonormal basis for  $\ell^2(\mathbf{G}) \otimes H$  is  $\{\delta_g \otimes u_\alpha : g \in \mathbf{G}, \alpha \in J\}$ .

We are interested in the following von Neumann algebra on  $\ell^2(\mathbf{G}) \otimes H$ :

$$\mathcal{M}_r(\mathbf{G}) \otimes \mathcal{B}(H).$$

By combining (1) and (2) in, we obtain the following trace function.

**Definition 6.3.** Let  $\{u_\alpha\}_{\alpha \in J}$  be an orthonormal basis for  $H$ . For every  $A \in (\mathcal{M}_r(\mathbf{G}) \otimes \mathcal{B}(H))^+$ , define

$$\mathrm{Tr}_{\mathbf{G}}(A) = (\tau \otimes \mathrm{Tr})(A) = \sum_{\alpha \in J} \langle A(\delta_h \otimes u_\alpha), \delta_h \otimes u_\alpha \rangle,$$

for  $h \in \mathbf{G}$ . This is independent of the choice of  $h$ , so it is convenient to use the  $h = \mathrm{id}$ , the group identity. It can be checked that this is a trace function according to Definition 3.1.

**Definition 6.4.** A projective Hilbert  $\mathbf{G}$ -module is a Hilbert  $\mathbf{G}$ -module  $V$  which is unitarily isomorphic to a closed submodule of a free Hilbert  $\mathbf{G}$ -module, i.e. a closed  $\mathbf{G}$ -invariant subspace in *some*  $\ell^2(\mathbf{G}) \otimes H$ .

Note that the embedding of  $V$  into  $\ell^2(\mathbf{G}) \otimes H$  is *not* part of the structure; only its existence is required.

For the moment, we will fix such an embedding. Let  $P_V$  denote the orthogonal projection from  $\ell^2(\mathbf{G}) \otimes H$  onto  $V$ . Note that  $P_V \in \mathcal{M}_r(\mathbf{G}) \otimes \mathcal{B}(H)$  because it commutes with all  $L_g \otimes I$ .

**Definition 6.5.** Let  $V$  be a projective Hilbert  $\mathbf{G}$ -module. Fix an embedding into a free Hilbert  $\mathbf{G}$ -module  $\ell^2(\mathbf{G}) \otimes H$ , and let  $P_V \in \mathcal{M}_r(\mathbf{G}) \otimes \mathcal{B}(H)$  denote the orthogonal projection onto  $V$ . The  $\mathbf{G}$ -dimension of  $V$  is defined as

$$\dim_{\mathbf{G}}(V) = \mathrm{Tr}_{\mathbf{G}}(P_V).$$

The amazing (but elementary) fact is that  $\dim_{\mathbf{G}}(V)$  does *not* depend on the choice of the embedding of  $V$  into a free Hilbert  $\mathbf{G}$ -module. Therefore it is a well-defined invariant of the projective Hilbert  $\mathbf{G}$ -module  $V$ . Furthermore, it satisfies the following properties:

- (i)  $\dim_{\mathbf{G}}(V) = 0 \iff V = \{0\}$ .

- (ii)  $\dim_{\mathbf{G}}(\ell^2(\mathbf{G})) = 1$ .
- (iii)  $\dim_{\mathbf{G}}(\ell^2(\mathbf{G}) \otimes H) = \dim_{\mathbb{C}}(H)$ .
- (iv)  $\dim_{\mathbf{G}}(V_1 \oplus V_2) = \dim_{\mathbf{G}}(V_1) + \dim_{\mathbf{G}}(V_2)$ .
- (v)  $V_1 \subseteq V_2 \Rightarrow \dim_{\mathbf{G}}(V_1) \leq \dim_{\mathbf{G}}(V_2)$ . Equality holds here if and only if  $V_1 = V_2$ .
- (vi) If  $0 \rightarrow U \rightarrow V \rightarrow W \rightarrow 0$  is a short *weakly exact sequence* of projective Hilbert  $\mathbf{G}$ -modules, then  $\dim_{\mathbf{G}}(V) = \dim_{\mathbf{G}}(U) + \dim_{\mathbf{G}}(W)$ .
- (vii)  $V$  and  $W$  are *weakly isomorphic*, then  $\dim_G(V) = \dim_G(W)$ .

We call a sequence of  $U \xrightarrow{i} V \xrightarrow{p} W$  of projective Hilbert  $\mathbf{G}$ -modules *weakly exact* at  $V$  if  $\text{Kernel}(p) = \text{cl}(\text{Image}(i))$ . A map of projective Hilbert  $\mathbf{G}$ -modules  $V \rightarrow W$  is a *weak isomorphism* if it is injective and has dense image.

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