

# Lecture 11

Intro to 2<sup>nd</sup> order ODEs

## 2<sup>nd</sup> Order ODEs:

$$\frac{d^2 y}{dt^2} = F(t, y, \frac{dy}{dt})$$

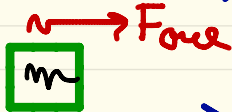
Initial conditions:  $\begin{cases} y(t_0) = y_0 \\ y'(t_0) = y'_0 \end{cases}$

Most important example:

Newton's 2<sup>nd</sup> Law of Motion:

mass  $\times$  acceleration = Force

$$m \frac{d^2 x}{dt^2} = F(t, x, \frac{dx}{dt})$$



Note: position and velocity at time  $t_0$  uniquely determine  $x(t)$ .

## Special Cases:

i)  $\frac{d^2 y}{dt^2} = F(t, \frac{dy}{dt})$

Let  $u = \frac{dy}{dt}$ :

$$\begin{cases} \frac{du}{dt} = F(t, u) & \text{1<sup>st</sup> order} \\ u(t_0) = u_0 \end{cases}$$

$$\Rightarrow u = U(t)$$

Then  $\begin{cases} \frac{dy}{dt} = U(t) \\ y(t_0) = y_0 \end{cases}$

$$\Rightarrow y(t) = y_0 + \int_{t_0}^t U(u) du$$

Example: Free fall with

air resistance:

$$m \frac{d^2 y}{dt^2} = -mg - \gamma \frac{dy}{dt}$$

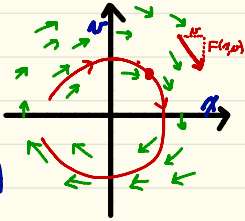
A diagram showing a green circle labeled 'm' representing a mass. A vertical y-axis points upwards. Two force vectors are shown: a green arrow pointing up labeled  $-\gamma \frac{dy}{dt} = F_{\text{drag}}$  and a black arrow pointing down labeled  $-mg = F_{\text{gravity}}$ . The mass is positioned between these two force vectors.

(2) Autonomous 2<sup>nd</sup> order ODE:

$$\frac{d^2x}{dt^2} = F(x, \frac{dx}{dt}) \Leftrightarrow \begin{cases} \frac{dx}{dt} = v \\ \frac{dv}{dt} = F(x, v) \end{cases}$$

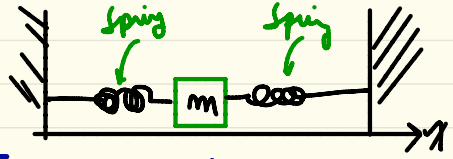
Phase plane

Solution  $\begin{cases} x = x(t) \\ v = v(t) \end{cases}$



Key idea: View  $v$  as a function of  $x$

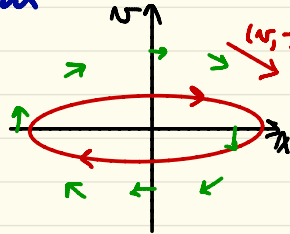
Example (Harmonic Oscillator)



Force =  $-kx$

$$m \frac{d^2x}{dt^2} = -kx$$

$$\Leftrightarrow \begin{cases} \frac{dx}{dt} = v \\ \frac{dv}{dt} = -kx/m \end{cases}$$



Conservation of Energy

$$m \frac{d^2x}{dt^2} = -kx$$

$$\Leftrightarrow m \frac{dv}{dt} = -kx$$

Chain Rule:  $\frac{dv}{dt} = \frac{dv}{dx} \frac{dx}{dt} = v \frac{dv}{dx}$

So

$$m v \frac{dv}{dx} = -kx$$

Integrate with respect to  $x$ :

$$\int m v \frac{dv}{dx} dx = - \int kx dx$$

$$\Rightarrow \frac{1}{2} m v^2 = -\frac{1}{2} kx^2 + E$$

$$\Rightarrow \frac{1}{2} m v^2 + \frac{1}{2} kx^2 = E$$

Constant of integration

kinetic + potential = constant.  
energy energy

Check:  $\frac{d}{dt} \left( \frac{1}{2} m v^2 + \frac{1}{2} k x^2 \right) = m v \frac{dv}{dt} + k x \frac{dx}{dt} = (-kx)v + (kx)v = 0$  ✓

Can use conservation of energy to solve ODE

$$\left\{ \begin{aligned} \frac{1}{2} m v^2 + \frac{1}{2} k x^2 &= E \text{ (constant)} \\ v &= \frac{dx}{dt} \end{aligned} \right.$$

$$\text{So } \frac{1}{2} m \left( \frac{dx}{dt} \right)^2 + \frac{1}{2} k x^2 = E$$

$$\Rightarrow \left( \frac{dx}{dt} \right)^2 = \frac{2E - kx^2}{m}$$

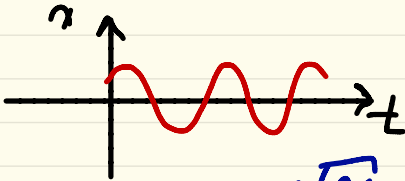
$$\left( \frac{dx}{dt} \right)^2 = \frac{k}{m} \left( \frac{2E}{k} - x^2 \right)$$

$$\frac{dx}{dt} = \pm \sqrt{\frac{k}{m}} \sqrt{\frac{2E}{k} - x^2}$$

$$\Rightarrow \int \frac{dx}{\sqrt{\frac{2E}{k} - x^2}} = \pm \sqrt{\frac{k}{m}} t + C$$

$$\Rightarrow \sin^{-1} \left( \frac{x}{\sqrt{\frac{2E}{k}}} \right) = \pm \sqrt{\frac{k}{m}} t + C$$

$$\text{Or } x(t) = \sqrt{\frac{2E}{k}} \sin \left( \sqrt{\frac{k}{m}} t - \beta \right)$$



$$\underline{\underline{Q}} \quad x(t) = \sqrt{\frac{2E}{k}} \cos \left( \sqrt{\frac{k}{m}} t - \varphi \right)$$

$$\varphi = \beta + \frac{\pi}{2}$$

$$\sin \alpha$$

$$= \cos \left( \alpha - \frac{\pi}{2} \right)$$

## Linear ODEs

$$\frac{dy}{dt} + p(t)y = \begin{cases} 0 & \text{(Homogeneous)} \\ f(t) & \text{(nonhomogeneous)} \end{cases}$$

$$\frac{d^2y}{dt^2} + p(t)\frac{dy}{dt} + q(t)y = \begin{cases} 0 & \text{(Homogeneous)} \\ f(t) & \text{(nonhomogeneous)} \end{cases}$$

Shorthand notation:

$$L[y] = y'' + p(t)y' + q(t)y$$

$L$  is an example  
of a linear

operator:

$$\begin{aligned} L[C_1 y_1(t) + C_2 y_2(t)] \\ = C_1 L[y_1(t)] + C_2 L[y_2(t)] \end{aligned}$$

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Example If  $L[y] = y'' + 4y$

Then  $L[\sin(2t)]$

$$\begin{aligned} &= (\sin(2t))'' + 4(\sin(2t)) \\ &= -4\sin(2t) + 4\sin(2t) = 0 \end{aligned}$$

So  $y(t) = \sin(2t)$  is a solution  
of the ODE

$$L[y] = y'' + 4y = 0$$

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Theorem Suppose  $p(t)$ ,  $q(t)$ , and  $f(t)$  are continuous on the interval  $a < t < b$ . If  $a < t_0 < b$ , then the initial value problem

$$\begin{aligned} L[y] &= y'' + p(t)y' + q(t)y = f(t) \\ y(t_0) &= y_0 \quad y'(t_0) = y'_0 \end{aligned}$$

Has a unique solution,  $y(t)$  defined for all  $a < t < b$ .

Bad news: Theorem

does not give a way to find the solution!

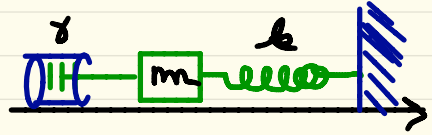
Good news: If  $p(t)=b$  and  $q(t)=c$  (i.e. constants), so

$$L[y] = y'' + by' + cy$$

Then there are methods for solving the I.V.P.

Main Example:

The (damped) harmonic oscillator:



$$m \frac{d^2x}{dt^2} = -kx - \gamma \frac{dx}{dt}$$

$$m \frac{d^2x}{dt^2} + \gamma \frac{dx}{dt} + kx = 0$$

Plan for the next 2 weeks:

(1) Study the homogeneous case:

$$L[y] = ay'' + by' + cy = 0$$

(2) Study the nonhomogeneous case

$$L[y] = f(t)$$

in special cases:

Example. Solve the I.V.P.

$$L[y] = y'' + 3y' + 2y = 0 \quad y(0) = 1, y'(0) = 2$$

Solution:

$$L[e^{rt}] = (r^2 + 3r + 2)e^{rt} = 0$$

$$\Leftrightarrow r^2 + 3r + 2 = 0$$

$$\Leftrightarrow (r+2)(r+1) = 0$$

$$\Rightarrow r = -2 \text{ or } r = -1$$

$$\text{So } y(t) = C_1 e^{-t} + C_2 e^{-2t}$$

is a solution for any  $C_1, C_2$

$$\begin{cases} y(0) = C_1 + C_2 = 1 \\ y'(0) = -C_1 - 2C_2 = 2 \end{cases}$$

$$\Rightarrow \begin{cases} C_1 = 4 \\ C_2 = -3 \end{cases}$$

$$\Rightarrow y(t) = 4e^{-t} - 3e^{-2t}$$

